# MA-207 Differential Equations II 

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## Gamma function

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For example, $\quad \Gamma\left(-\frac{5}{2}\right)=\frac{\Gamma\left(-\frac{3}{2}\right)}{-\frac{5}{2}}=\frac{\Gamma\left(-\frac{1}{2}\right)}{\left(-\frac{5}{2}\right)\left(-\frac{3}{2}\right)}=\frac{\Gamma\left(\frac{1}{2}\right)(=\sqrt{\pi})}{\left(-\frac{5}{2}\right)\left(-\frac{3}{2}\right)\left(-\frac{1}{2}\right)}$

Further

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\lim _{p \rightarrow 0} \Gamma(p)=\lim _{p \rightarrow 0} \frac{\Gamma(p+1)}{p}= \pm \infty
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The graph of Gamma function is shown below.

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\begin{aligned}
\Gamma(1 / 2) & =\int_{0}^{\infty} t^{-1 / 2} e^{-t} d t \\
& =2 \int_{0}^{\infty} e^{-s^{2}} d s \quad \text { (use the substitution } t=s^{2} \text { ) } \\
& =\sqrt{\pi}
\end{aligned}
$$

## By translating,

$$
\begin{aligned}
& \Gamma(1 / 2) \quad=\sqrt{\pi} \quad \approx 1.772 \\
& \Gamma(-1 / 2)=\frac{\Gamma(1 / 2)}{-1 / 2}=-2 \sqrt{\pi} \approx-3.545 \\
& \Gamma(-3 / 2)=\frac{\Gamma(-1 / 2)}{-3 / 2}=\frac{4}{3} \sqrt{\pi} \quad \approx 2.363 \\
& \Gamma(3 / 2) \quad=\frac{1}{2} \Gamma(1 / 2) \quad=\frac{1}{2} \sqrt{\pi} \quad \approx 0.886 \\
& \Gamma(5 / 2) \quad=\frac{3}{2} \Gamma(3 / 2) \quad=\frac{3}{4} \sqrt{\pi} \quad \approx 1.329 \\
& \Gamma(7 / 2)=\frac{5}{2} \Gamma(5 / 2)=\frac{15}{8} \sqrt{\pi} \approx 3.323
\end{aligned}
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Bessel functions have applications in physics and engineering:
Since $x=0$ is a regular singular point of $(*)$, we get a Frobenius solution, called Bessel function of first kind.
The second linearly independent solution of $(*)$ is called Bessel function of second kind.
For Frobenius solution, put $y=x^{r} \sum_{n=0}^{\infty} a_{n}(r) x^{n} \quad a_{0}=1$.

The indicial equation, i.e. coefficient of $x^{r}$, for Bessel equation $x^{2} y^{\prime \prime}+x y^{\prime}+\left(x^{2}-p^{2}\right) y=0 \quad$ is

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I(r)=r(r-1)+r-p^{2}=r^{2}-p^{2}=0
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The roots are $r_{1}=p$ and $r_{2}=-p$.
For recurrence relations, equating coefficient of $x^{n+r}$ to 0 (for $n \geq 1$ ) we get

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\left[(r+n)^{2}-p^{2}\right] a_{n}(r)+a_{n-2}(r)=0 \quad n \geq 2
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& {\left[(r+n)^{2}-p^{2}\right] a_{n}(r)+a_{n-2}(r)=0 \quad n \geq 2} \\
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So all odd terms $a_{2 n+1}(r)=0$.

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\begin{aligned}
a_{2 n}(r) & =\frac{-1}{(r+2 n)^{2}-p^{2}} a_{2 n-2} \\
& =\frac{(-1)^{n}}{\left((r+2)^{2}-p^{2}\right)\left((r+4)^{2}-p^{2}\right) \ldots\left((r+2 n)^{2}-p^{2}\right)}
\end{aligned}
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For Frobenius solution, set $r=p$ the larger root.

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& =\frac{(-1)^{n}}{(2(2 p+2))(4(2 p+4)) \ldots(2 n(2 p+2 n))} \\
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The solution

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y_{1}(x)=x^{p} \sum_{n \geq 0} \frac{(-1)^{n}}{2^{2 n} n!(1+p) \ldots(n+p)} x^{2 n}
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converges on $(0, \infty)$.

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converges on $(0, \infty)$.
Multiply $y_{1}(x)$ by $\frac{1}{2^{p} \Gamma(1+p)}$

$$
J_{p}(x):=\left(\frac{x}{2}\right)^{p} \sum_{n \geq 0} \frac{(-1)^{n}}{n!\Gamma(n+p+1)}\left(\frac{x}{2}\right)^{2 n} \quad x>0 .
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This is called the Bessel function of first kind of order $p$.

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\begin{aligned}
J_{0}(x) & =\sum_{n \geq 0} \frac{(-1)^{n}}{n!n!}\left(\frac{x}{2}\right)^{2 n} \\
& =1-\left(\frac{x}{2}\right)^{2}+\frac{1}{2!2!}\left(\frac{x}{2}\right)^{4}-\frac{1}{3!3!}\left(\frac{x}{2}\right)^{6}+\ldots
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The Bessel function of order 1 is

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& =\frac{x}{2}-\frac{1}{1!2!}\left(\frac{x}{2}\right)^{3}+\frac{1}{2!3!}\left(\frac{x}{2}\right)^{5}+\ldots
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Further, they satisfy derivative identities similar to $\cos x$ and $\sin x$.

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J_{0}^{\prime}(x)=-J_{1}(x) \quad\left[x J_{1}(x)\right]^{\prime}=x J_{0}(x)
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## Second independent solution of Bessel equation

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The analysis to get a second independent solution of the Bessel equation splits into the following cases

- $2 p$ is not an integer
- $2 p$ is an odd positive integer
- $2 p$ is an even positive integer
- $p=0$


## Second independent solution of Bessel equation

Case 1: $2 p$ is not an integer.

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for $r=-p$, we obtain

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Multiplying by $\frac{1}{2^{-p} \Gamma(1-p)}$

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It is unbounded near $x=0$.

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Recall that the second solution is given by

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y_{2}(x)=\sum_{n \geq 0} A_{n}^{\prime}(-p) x^{n-p}+\sum_{n \geq 0} A_{n}(-p) x^{n-p} \log x
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Multiplying by $\frac{1}{2^{-p} \Gamma(1-p)}$

$$
J_{-p}(x):=\left(\frac{x}{2}\right)^{-p} \sum_{n \geq 0} \frac{(-1)^{n}}{n!\Gamma(n-p+1)}\left(\frac{x}{2}\right)^{2 n} \quad x>0 .
$$

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As before, $A_{2 n+1}(r)=0$.
The polynomial $\prod_{i=1}^{n}\left((r+2 i)^{2}-p^{2}\right)$ evaluated at $r=-p$, is
$\prod_{i=1}^{n} 4 i(i-p)$,
Thus, if $n<p$, then $a_{2 n}(r)$ is analytic in a neighborhood of $-p$.
Thus, if $n<p$, then $A_{2 n}(-p)=0$ and

$$
A_{2 n}^{\prime}(-p)=a_{2 n}(-p)=\frac{(-1)^{n}}{2^{2 n} n!(1-p) \ldots(n-p)}=\frac{1}{2^{2 n} n!(p-n)!}
$$

If $n \geq p$, then

$$
\begin{aligned}
A_{2 n}(-p) & =\frac{2(-1)^{n}}{2^{2 n} n!(1-p) \ldots(-1) \cdot 1 \cdot 2 \cdots(n-p)} \\
& =\frac{-2(-1)^{n-p}}{2^{2 n} n!(p-1)!(n-p)!}
\end{aligned}
$$

Define
$f(r):=\left(\prod_{i=1}^{p-1}\left((r+2 i)^{2}-p^{2}\right)\right)(r+3 p)\left(\prod_{i=p+1}^{n}\left((r+2 i)^{2}-p^{2}\right)\right)$

Then

$$
A_{2 n}(r) f(r)=(-1)^{n}
$$

Differentiating the above and setting $r=-p$ we get

$$
A_{2 n}^{\prime}(-p) f(-p)+A_{2 n}(-p) f^{\prime}(-p)=0
$$

Taking log and differentiating $(*)$ we get

$$
\begin{gathered}
f^{\prime}(-p)=f(-p)\left(\frac{1}{2 p}+\sum_{i \in\{1,2, \ldots, n\} \backslash p} \frac{1}{2 i}+\frac{1}{2(i-p)}\right) \\
=f(-p)\left(\frac{H_{n}}{2}-\frac{H_{p-1}}{2}+\frac{H_{n-p}}{2}\right), \\
H_{0}=0, \quad H_{n}=1+\frac{1}{2}+\cdots+\frac{1}{n}
\end{gathered}
$$

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Thus,

$$
\begin{aligned}
A_{2 n}^{\prime}(-p) & =-A_{2 n}(-p)\left(\frac{H_{n}}{2}-\frac{H_{p-1}}{2}+\frac{H_{n-p}}{2}\right) \\
& =\frac{2(-1)^{n-p}}{2^{2 n} n!(p-1)!(n-p)!}\left(\frac{H_{n}}{2}-\frac{H_{p-1}}{2}+\frac{H_{n-p}}{2}\right)
\end{aligned}
$$

Thus, we get

$$
\begin{aligned}
& y_{2}(x)=\sum_{n=0}^{p-1} \frac{1}{2^{2 n} n!(p-n)!} x^{2 n-p}+ \\
& \sum_{n \geq p} \frac{(-1)^{n-p}}{2^{2 n} n!(p-1)!(n-p)!}\left(H_{n}-H_{p-1}+H_{n-p}\right) x^{2 n-p}+ \\
& \quad-\sum_{n \geq p} \frac{2(-1)^{n-p}}{2^{2 n} n!(p-1)!(n-p)!} x^{2 n-p} \log x
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y_{2}(x)=J_{0}(x) \ln x-\sum_{n \geq 1} \frac{(-1)^{n} H_{n}}{2^{2 n}(n!)^{2}} x^{2 n} \quad x>0
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where $y_{1}(x)=J_{0}(x)=\sum_{n \geq 0} \frac{(-1)^{n}}{2^{2 n}(n!)^{2}} x^{2 n}$ is Frobenius solution.

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Both exhibit singular behavior at 0 . Near $0, J_{1 / 2}(x)$ is bounded but does not have a bounded derivative, while $J_{-1 / 2}(x)$ is unbounded near 0 .

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(1) The above is a well defined power series once we know that the Gamma function never vanishes.
(2) If $p \notin\{0,1,2, \ldots\} J_{p}(x)$ and $J_{-p}(x)$ are the two independent solutions of the Bessel equation.
(3) If $p \in\{0,1,2, \ldots\}$ then $J_{-p}(x)=(-1)^{p} J_{p}(x)$. Thus, in this case the second solution is not $J_{-p}(x)$.

## Bessel identities

$$
\begin{aligned}
& \text { (1) } \frac{d}{d x}\left[x^{p} J_{p}(x)\right]=x^{p} J_{p-1}(x) \\
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Add and subtract (3) and (4) to get (5) and (6).

## Consequences of Bessel identities

Problem: Let $p>0$. Show that between any two consecutive zeros of $J_{p}(x)$, there exists precisely one zero of $J_{p-1}(x)$ and precisely one zero of $J_{p+1}(x)$

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we get $J_{p-1}(b)=0$.
Repeating the above argument with the identity $\left[x^{-p} J_{p}(x)\right]^{\prime}=-x^{-p} J_{p+1}(x)$, we get that $J_{p+1}(x)$ has a root in $(c, d)$.

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Hence $a=2$ and $c=0$.

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For example,

- $J_{3 / 2}(x)=\frac{1}{x} J_{1 / 2}(x)-J_{-1 / 2}(x)$

$$
=\sqrt{\frac{2}{\pi x}}\left(\frac{\sin x}{x}-\cos x\right)
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to see that $J_{p}(x)$ are elementary functions for $p \in \mathbb{Z}+\frac{1}{2}$.
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- $J_{3 / 2}(x)=\frac{1}{x} J_{1 / 2}(x)-J_{-1 / 2}(x)$

$$
=\sqrt{\frac{2}{\pi x}}\left(\frac{\sin x}{x}-\cos x\right)
$$

- $J_{-3 / 2}(x)=-\frac{1}{x} J_{-1 / 2}(x)-J_{1 / 2}(x)$

$$
=-\sqrt{\frac{2}{\pi x}}\left(\frac{\cos x}{x}+\sin x\right)
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- $J_{\frac{5}{2}}(x)=\frac{3}{x} J_{\frac{3}{2}}(x)-J_{\frac{1}{2}}(x)$

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These functions are called spherical Bessel functions as they arise in solving wave equations in spherical coordinates.

## Theorem (Liouville)

$J_{m+\frac{1}{2}}(x)$ 's are the only elementary Bessel functions.

Remark. Integrating some of the Bessel identities we get

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& \frac{d}{d x}\left[x^{p} J_{p}(x)\right]=x^{p} J_{p-1}(x) \\
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## Qualitative properties of solutions

It is rarely possible to solve 2nd order linear ODE

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y^{\prime \prime}+P(x) y^{\prime}+Q(x) y=0
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in terms of familiar elementary functions.
Then how do we understand the nature and properties of solutions.

It is surprising that we can obtain quite a bit of information about the solution from the ODE itself.

## Theorem (Sturm separation theorem)

If $y_{1}(x)$ and $y_{2}(x)$ are linearly independent solns of

$$
y^{\prime \prime}+P(x) y^{\prime}+Q(x) y=0
$$

$P, Q$ continuous on $(a, b)$. Then
(1) $y_{1}(x)$ and $y_{2}(x)$ have no common zero in $(a, b)$.
(2) Between any two successive zeros of $y_{1}(x)$, there is exactly one zero of $y_{2}(x)$ and vice versa.

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In particular, since $y_{1}$ and $y_{2}$ are linearly independent, the Wronskian is nonzero and so it never vanishes. This proves (1).

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First let us show $y_{2}$ has a zero in $\left(x_{1}, x_{2}\right)$.
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If $y_{2}(x)$ had two zeros in the interval $x_{1}<\alpha<\beta<x_{2}$, then by the same reasoning, $y_{1}$ will have a zero in $(\alpha, \beta)$, which contradicts the assumption that $x_{1}$ and $x_{2}$ are successive zeros of $y_{1}$.

As a consequence, if $y_{1}$ and $y_{2}$ are linearly independent solution of $y^{\prime \prime}+P(x) y^{\prime}+Q(x) y=0, P, Q$ continuous on $(a, b)$ then the number of zeros of $y_{1}$ and $y_{2}$ on $(a, b)$ differ by at most 1 .

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- For further discussion, we need that any ODE in the "standard" form $y^{\prime \prime}+P(x) y^{\prime}+Q(x) y=0$ can be written in the "normal" form $u^{\prime \prime}+q(x) u=0$.

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It is clear that the zeros of $u$ are the same as those of $y$. Also note that we need $P(x)$ to be once differentiable.

## Theorem

Let $u(x)$ be a non-trivial solution of $u^{\prime \prime}+q(x) u=0$ on the interval $(\alpha, \beta)$, with $q(x)$ continuous. Let $[a, b] \subset(\alpha, \beta)$ be a finite interval. Then $u(x)$ has at most finite number of zeros in $[a, b]$. Hence if $u(x)$ has infinitely many zeros on $(0, \infty)$, then the set of zeros of $u(x)$ are not bounded.

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Proof. Assume $u(x)$ has infinitely many zeros in $[a, b]$. Then $\exists x_{0} \in[a, b]$ and a sequence of zeros $x_{n} \neq x_{0}$ such that $x_{n} \rightarrow x_{0}$ as $n \rightarrow \infty$.

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$u\left(x_{0}\right)=\lim _{x_{n} \rightarrow x_{0}} u\left(x_{n}\right)=0$ ( $u$ is continuous) and

$$
u^{\prime}\left(x_{0}\right)=\lim _{x_{n} \rightarrow x_{0}} \frac{u\left(x_{n}\right)-u\left(x_{0}\right)}{x_{n}-x_{0}}=0
$$

This contradicts the fact that the Wronskian at $x_{0}$ is nonzero.

Let $u(x)$ be a non-trivial solution of $u^{\prime \prime}+q(x) u=0$. If $q(x)<0$ in $(a, b)$ and continuous then $u(x)$ has atmost one zero in $(a, b)$.

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If necessary, multiply by -1 and assume that $u^{\prime}\left(x_{0}\right)>0$.

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Since $u^{\prime \prime}(x)=-q(x) u(x)>0$ on $\left(x_{0}, x_{1}\right), u^{\prime}(x)$ is an increasing function on ( $x_{0}, x_{1}$ ).

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But this is not possible as $u^{\prime}$ is increasing on $\left(x_{0}, x_{1}\right)$.

Let $u(x)$ be a non-trivial solution of $u^{\prime \prime}+q(x) u=0$ Let $q(x)$ be continuous and $q(x)>0$ for all $x>x_{0}>0$.
If $\int_{x_{0}}^{\infty} q(x) d x=\infty$,
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Proof. Assume $u(x)$ has only finitely many zeros on $(0, \infty)$.
Then there is $x_{1}>x_{0}$ such that $u(x) \neq 0$ for $x \geq x_{1}$. Assume $u(x)>0$ for $x \geq x_{1}$.

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If we show that $u^{\prime}\left(x_{2}\right)<0$ for some $x_{2}>x_{1}$, then we get for $x>x_{2}$

$$
\begin{aligned}
u(x) & =\int_{x_{2}}^{x} u^{\prime}(t) d t+u\left(x_{2}\right) \leq \int_{x_{2}}^{x} u^{\prime}\left(x_{2}\right) d t+u\left(x_{2}\right) \\
& \leq u^{\prime}\left(x_{2}\right)\left(x-x_{2}\right)+u\left(x_{2}\right)
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Thus, $u^{\prime}(x)=-u(x) v(x)$ and this shows that $u^{\prime}(x)<0$ for $x$ large.

## Theorem

In Bessel equation $x^{2} y^{\prime \prime}+x y^{\prime}+\left(x^{2}-p^{2}\right) y=0$ Substituting $u(x)=\sqrt{x} y(x)$, we get

$$
u^{\prime \prime}+\left[1+\frac{1-4 p^{2}}{4 x^{2}}\right] u=0
$$

$q(x)=1+\frac{1-4 p^{2}}{4 x^{2}}$ is continuous and $q(x)>0$ for $x>x_{0}>0$.
Further,

$$
\int_{x_{0}}^{\infty}\left(1+\frac{1-4 p^{2}}{4 x^{2}}\right) d x=\infty
$$

By previous theorem, $u(x)$, hence any Bessel function has infinitely many zeros on $(0, \infty)$.

## Corollary

Let $Z^{(p)}$ be the set of zeros of Bessel function $J_{p}(x)$ on $(0, \infty)$.

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Question. What is the limit of $x_{n+1}-x_{n}$ as $n \rightarrow \infty$ ?

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We will need the Sturm comparison theorem.

## Theorem (Sturm Comparison theorem)

Let $y(x)$ be a non-trivial solutions of

$$
y^{\prime \prime}+q(x) y=0
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and $z(x)$ be a non-trivial solutions of

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where $q(x)>r(x)>0$ are continuous.
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Zeros of $y(x)$ are $\pi / 2$ apart and that of $z(x)$ are $\pi$ apart.

## Proof of Sturm Comparison theorem.

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Integrating from $x_{1}$ to $x_{2}$, we get

$$
W\left(x_{2}\right)-W\left(x_{1}\right)>0 \Longrightarrow W\left(x_{2}\right)>W\left(x_{1}\right)
$$

But $W\left(x_{1}\right)=y\left(x_{1}\right) z^{\prime}\left(x_{1}\right)>0$ and $W\left(x_{2}\right)=y\left(x_{2}\right) z^{\prime}\left(x_{2}\right)<0$, a contradiction.

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Let $y_{p}(x)$ be a non-trivial solution of Bessel equation. Then we get

## Theorem

- $p<1 / 2 \Longrightarrow$ Between any two roots of $\alpha \cos x+\beta \sin x$ there is a root of $y_{p}(x)$.
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To see this, consider the function $f:=u(x-b+a)$ defined on the interval $(b, \infty)$.

It is a trivial check that $f$ satisfies the differential equation

$$
f^{\prime \prime}+r(x) f=0 \quad r(x):=q(x-b+a)
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Clearly,

- $b<x_{0} \Longrightarrow a<x_{0}-b+a$
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Next we claim that the difference between any two successive roots of $u$ is strictly less than $\pi$.

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But now we can choose $\alpha, \beta \in \mathbb{R}$ such that $\alpha \cos x+\beta \sin x$ has two roots in ( $a, b$ ), which contradicts Sturm's comparison theorem.

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Thus, we have proved that if $\left\{x_{n}\right\}$ are the roots of $u$ in increasing order, then the difference $x_{n+1}-x_{n}$ is strictly increasing and bounded above by $\pi$.

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Next let us show that these differences converge to $\pi$. If not, then $\left(x_{n+1}-x_{n}\right) \rightarrow \gamma<\pi$. Choose $1<\delta$, sufficiently close to 1 such that $\gamma<\frac{\pi}{\delta}<\pi$.

The function $q(x)$ is decreasing to 1 . Therefore, there is a $x_{0} \in \mathbb{R}$, sufficiently large, such that $q\left(x_{0}\right)<\delta^{2}$. Apply Sturm's comparison on the interval $\left(x_{0}, \infty\right)$ to the differential equations $u^{\prime \prime}+q(x) u=0$ and $z^{\prime \prime}+\delta^{2} z=0$.

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Thus, between any two roots of $u$ there is a root of $z$. Let $a$ and $b$ be two consecutive roots of $u$ such that $x_{0}<a<b$. Since $b-a<\gamma<\frac{\pi}{\delta}$, find $a^{\prime}$ and $b^{\prime}$ such that $x_{0}<a^{\prime}<a<b<b^{\prime}$ and $b^{\prime}-a^{\prime}=\frac{\pi}{\delta}$.

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Find $\alpha$ and $\beta$ such that the function $\alpha \cos \delta x+\beta \sin \delta x$ vanishes at $a^{\prime}$. This function is a solution to the ODE $z^{\prime \prime}+\delta^{2} z=0$. The next root of this function is at $a^{\prime}+\frac{\pi}{\delta}=b^{\prime}$. Thus, we get a contradiction to Sturm's theorem which says that there is a root of this function in the interval $(a, b)$.

Thus, we have proved

## Theorem

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Similarly, we can prove that if $p>1 / 2$ then the sequence of difference of roots of $u$ is decreasing and tends to $\pi$.


The first few zeroes of Bessel functions are tabulated below.

|  | $J_{0}(x)$ | $J_{1}(x)$ | $J_{2}(x)$ | $J_{3}(x)$ | $J_{4}(x)$ | $J_{5}(x)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 2.4048 | 3.8317 | 5.1356 | 6.3802 | 7.5883 | 8.7715 |
| 2 | 5.5201 | 7.0156 | 8.4172 | 9.7610 | 11.0647 | 12.3386 |
| 3 | 8.6537 | 10.1735 | 11.6198 | 13.0152 | 14.3725 | 15.7002 |
| 4 | 11.7915 | 13.3237 | 14.7960 | 16.2235 | 17.6160 | 18.9801 |
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Question. Why are we concerned with zeros of Bessel function $J_{p}(x)$ ?
It is often required in mathematical physics to expand a given function in terms of Bessel functions.

Simplest and most useful expansions are of the form

$$
f(x)=\sum_{n=1}^{\infty} a_{n} J_{p}\left(\lambda_{p, n} x\right)=a_{1} J_{p}\left(\lambda_{p, 1} x\right)+a_{2} J_{p}\left(\lambda_{p, 2} x\right)+\ldots
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Qn. How to compute the coefficients $a_{n}$ ?
Remark: For a scalar $a$, the scaled Bessel functions $J_{p}(a x)$ are solutions of

$$
x^{2} y^{\prime \prime}+x y^{\prime}+\left(a^{2} x^{2}-p^{2}\right) y=0
$$

known as scaled Bessel equation.

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This is similar to the previous inner product except that $f(x) g(x)$ is now multiplied by $x$ and the interval of integration is from 0 to 1 .

We call a function on $[0,1]$ square integrable with respect to this inner product if

$$
\int_{0}^{1} x f(x)^{2} d x<\infty
$$

The multiplying factor $x$ is called a weight function.

Fix $p \geq 0$. Let $Z^{(p)}=\left\{\lambda_{p, 1}, \lambda_{p, 2}, \ldots\right\}$ denote the set of zeros of $J_{p}(x)$ on $(0, \infty)$.

## Theorem

The set of scaled Bessel functions

$$
\left\{J_{p}\left(\lambda_{p, 1} x\right), J_{p}\left(\lambda_{p, 2} x\right), \ldots\right\}
$$

form an orthogonal family w.r.t. above inner product, i.e. $\left\langle J_{p}\left(\lambda_{p, k} x\right), J_{p}\left(\lambda_{p, l} x\right)\right\rangle:=$

$$
\int_{0}^{1} x J_{p}\left(\lambda_{p, k} x\right) J_{p}\left(\lambda_{p, l} x\right) d x= \begin{cases}\frac{1}{2}\left[J_{p+1}\left(\lambda_{p, k}\right)\right]^{2} & \text { if } k=l \\ 0 & \text { if } k \neq l\end{cases}
$$

## Theorem

Fix $p \geq 0$ and $Z^{(p)}=\left\{\lambda_{p, 1}, \lambda_{p, 2}, \ldots\right\}$ : zeros of $J_{p}(x)$ on $(0, \infty)$. Any square-integrable function $f(x)$ on $[0,1]$ can be expanded in a series of scaled Bessel functions $J_{p}\left(\lambda_{p, n} x\right)$ as

$$
f(x)=\sum_{n \geq 1} c_{n} J_{p}\left(\lambda_{p, n} x\right)
$$

where

$$
c_{n}=\frac{2}{\left[J_{p+1}\left(\lambda_{p, n}\right)\right]^{2}} \int_{0}^{1} x f(x) J_{p}\left(\lambda_{p, n} x\right) d x
$$

This is Fourier-Bessel series of $f(x)$ for parameter $p$.

Example. Let us compute the Fourier-Bessel series (for $p=0$ ) of $f(x)=1$ in the interval $0 \leq x \leq 1$.

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\int_{0}^{1} x J_{0}\left(\lambda_{0, n} x\right) d x=\left.\frac{1}{\lambda_{0, n}} x J_{1}\left(\lambda_{0, n} x\right)\right|_{0} ^{1}=\frac{J_{1}\left(\lambda_{0, n}\right)}{\lambda_{0, n}}
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Thus, the Fourier-Bessel series of $f(x)$ is

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By next theorem, this converges to 1 for $0<x<1$.

## Convergence in norm

Fourier-Bessel series converges to $f(x)$ in norm, i.e.

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\left\|f(x)-\sum_{n=1}^{m} c_{n} J_{p}\left(\lambda_{p, n} x\right)\right\| \text { converges to } 0 \text { as } m \rightarrow \infty
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For pointwise convergence, we have
Bessel expansion theorem
Assume $f$ and $f^{\prime}$ have at most a finite number of jump discontinuities in $[0,1]$, then the Bessel series converges for $0<x<1$ to

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at $x=0$, if $p>0$ then it converges to 0 .

## Proof of orthogonality of scaled Bessel functions

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\begin{aligned}
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Multiply by $v$ and $u$ resp. and subtract, we get

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\left(v u^{\prime \prime}-u v^{\prime \prime}\right)+\frac{1}{x}\left(v u^{\prime}-u v^{\prime}\right)+\left(a^{2}-b^{2}\right) u v=0
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& \left(u^{\prime} v-v^{\prime} u\right)^{\prime}+\frac{1}{x}\left(u^{\prime} v-v^{\prime} u\right)=\left(b^{2}-a^{2}\right) u v \\
& \left(x\left(u^{\prime} v-v^{\prime} u\right)\right)^{\prime}=\left(b^{2}-a^{2}\right) x u v
\end{aligned}
$$

$$
\left(b^{2}-a^{2}\right) \int_{0}^{1} x u v d x=\left.\left[x\left(u^{\prime} v-v^{\prime} u\right)\right]\right|_{0} ^{1}=\left(u^{\prime} v-v^{\prime} u\right)(1)
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So if $a=\lambda_{p, k}$ and $b=\lambda_{p, l}$ are distinct, then

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\int_{0}^{1} x J_{p}\left(\lambda_{p, k} x\right) J_{p}\left(\lambda_{p, l} x\right) d x=0
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To compute the norm of $J_{p}\left(\lambda_{p, k} x\right)$, consider

$$
\begin{aligned}
& 2 x^{2} u^{\prime}\left[u^{\prime \prime}+\frac{1}{x} u^{\prime}+\left(a^{2}-\frac{p^{2}}{x^{2}}\right) u\right]=0 \\
& \quad=\left[x^{2} u^{\prime 2}+\left(a^{2} x^{2}-p^{2}\right) u^{2}\right]^{\prime}-2 a^{2} x u^{2}
\end{aligned}
$$

Integrate on $[0,1]$,

$$
2 a^{2} \int_{0}^{1} x u^{2} d x=\left.\left[x^{2} u^{\prime 2}+\left(a^{2} x^{2}-p^{2}\right) u^{2}\right]\right|_{0} ^{1}
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Since $p \geq 0,(p u(0))^{2}=\left(p J_{p}(0)\right)^{2}=0$.
Thus, $\left(x^{2} u^{\prime 2}+\left(a^{2} x^{2}-p^{2}\right) u^{2}\right)(0)=0$.

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$$

Put $a=\lambda_{p, k}$ to get

$$
2 \lambda_{p, k}^{2} \int_{0}^{1} x J_{p}\left(\lambda_{p, k} x\right)^{2} d x=\lambda_{p, k}^{2} J_{p}^{\prime}\left(\lambda_{p, k}\right)^{2}
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Thus,

$$
\int_{0}^{1} x J_{p}\left(\lambda_{p, k} x\right)^{2} d x=\frac{1}{2} J_{p}^{\prime}\left(\lambda_{p, k}\right)^{2}=\frac{1}{2} J_{p+1}\left(\lambda_{p, k}\right)^{2}
$$

for last equality, use $x=\lambda_{p, k}$ in $J_{p}^{\prime}(x)-\frac{p}{x} J_{p}(x)=J_{p+1}(x)$

