MA-207 Differential Equations II

Ronnie Sebastian



Department of Mathematics Indian Institute of Technology Bombay Powai, Mumbai - 76

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For any real number $p \ge 1$,

$$\Gamma(p+1) = \lim_{x \to \infty} \int_0^x t^p e^{-t} dt = p \left(\lim_{x \to \infty} \int_0^x t^{p-1} e^{-t} dt \right)$$

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$$\Gamma(p+1) = p \Gamma(p) \implies \Gamma(p) = \frac{\Gamma(p+1)}{p}$$
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For example,
$$\Gamma(-\frac{5}{2}) = \frac{\Gamma(-\frac{3}{2})}{-\frac{5}{2}} = \frac{\Gamma(-\frac{1}{2})}{(-\frac{5}{2})(-\frac{3}{2})} = \frac{\Gamma(\frac{1}{2})(=\sqrt{\pi})}{(-\frac{5}{2})(-\frac{3}{2})(-\frac{1}{2})}$$

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$$\lim_{p\to 0} \Gamma(p) = \lim_{p\to 0} \frac{\Gamma(p+1)}{p} = \pm \infty$$

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according as $p \to 0$ from right or left. The graph of Gamma function is shown below.



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$$\frac{1}{\Gamma(p)} = \frac{p}{\Gamma(p+1)}$$

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$$\begin{split} \Gamma(1/2) &= \int_0^\infty t^{-1/2} e^{-t} dt \\ &= 2 \int_0^\infty e^{-s^2} ds \quad \text{(use the substitution } t = s^2\text{)} \\ &= \sqrt{\pi} \end{split}$$

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By translating,

$$\begin{split} \Gamma(1/2) &= \sqrt{\pi} &\approx 1.772 \\ \Gamma(-1/2) &= \frac{\Gamma(1/2)}{-1/2} &= -2\sqrt{\pi} &\approx -3.545 \\ \Gamma(-3/2) &= \frac{\Gamma(-1/2)}{-3/2} &= \frac{4}{3}\sqrt{\pi} &\approx 2.363 \\ \Gamma(3/2) &= \frac{1}{2}\Gamma(1/2) &= \frac{1}{2}\sqrt{\pi} &\approx 0.886 \\ \Gamma(5/2) &= \frac{3}{2}\Gamma(3/2) &= \frac{3}{4}\sqrt{\pi} &\approx 1.329 \\ \Gamma(7/2) &= \frac{5}{2}\Gamma(5/2) &= \frac{15}{8}\sqrt{\pi} &\approx 3.323 \end{split}$$

Bessel functions

$$x^{2}y'' + xy' + (x^{2} - p^{2})y = 0 \qquad p \ge 0 \qquad (*)$$

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For Frobenius solution, put
$$y = x^r \sum_{n=0}^{\infty} a_n(r) x^n$$
 $a_0 = 1$.

$$I(r) = r(r-1) + r - p^{2} = r^{2} - p^{2} = 0$$

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The roots are $r_1 = p$ and $r_2 = -p$. For recurrence relations, equating coefficient of x^{n+r} to 0 (for $n \ge 1$) we get

 $[(r+n)^2 - p^2]a_n(r) + a_{n-2}(r) = 0 \quad n \ge 2$

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So all odd terms $a_{2n+1}(r) = 0$.

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So all odd terms $a_{2n+1}(r) = 0$.

$$a_{2n}(r) = \frac{-1}{(r+2n)^2 - p^2} a_{2n-2}$$
$$= \frac{(-1)^n}{((r+2)^2 - p^2)((r+4)^2 - p^2)\dots((r+2n)^2 - p^2)}$$

For Frobenius solution, set r = p the larger root.

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$$= \frac{(-1)^n}{(2(2p+2))(4(2p+4))\dots(2n(2p+2n))}$$
$$= \frac{(-1)^n}{2^{2n}n!(1+p)\dots(n+p)}$$

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The solution $y_1(x) = x^p \sum_{n \ge 0} \frac{(-1)^n}{2^{2n}n!(1+p) \dots (n+p)} x^{2n}$
converges on $(0, \infty)$.
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converges on $(0, \infty)$.
Multiply $y_1(x)$ by $\frac{1}{2^p \Gamma(1+p)}$
 $J_p(x) := \left(\frac{x}{2}\right)^p \sum_{n \ge 0} \frac{(-1)^n}{n! \Gamma(n+p+1)} \left(\frac{x}{2}\right)^{2n} \quad x > 0.$

This is called the Bessel function of first kind of order p_{\pm} , z_{\pm}

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The Bessel function of order 0 is

$$J_0(x) = \sum_{n \ge 0} \frac{(-1)^n}{n! n!} \left(\frac{x}{2}\right)^{2n}$$

= $1 - \left(\frac{x}{2}\right)^2 + \frac{1}{2! 2!} \left(\frac{x}{2}\right)^4 - \frac{1}{3! 3!} \left(\frac{x}{2}\right)^6 + \dots$

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$$J_1(x) = \sum_{n \ge 0} \frac{(-1)^n}{n!(n+1)!} \left(\frac{x}{2}\right)^{2n+1}$$
$$= \frac{x}{2} - \frac{1}{1!2!} \left(\frac{x}{2}\right)^3 + \frac{1}{2!3!} \left(\frac{x}{2}\right)^5 + \dots$$

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Further, they satisfy derivative identities similar to $\cos x$ and $\sin x$.

$$J_0'(x) = -J_1(x) \qquad [xJ_1(x)]' = xJ_0(x)$$

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Recall $r_1 = p$ and $r_2 = -p$ are roots of indicial equation. So that $r_1 - r_2 = 2p$. Recall $r_1 = p$ and $r_2 = -p$ are roots of indicial equation. So that $r_1 - r_2 = 2p$.

The analysis to get a second independent solution of the Bessel equation splits into the following cases

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The analysis to get a second independent solution of the Bessel equation splits into the following cases

- 2p is not an integer
- 2p is an odd positive integer
- 2p is an even positive integer

•
$$p = 0$$

Case 1: 2p is not an integer.

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Solving the recursion

$$[(r+n)^2 - p^2]a_n(r) + a_{n-2}(r) = 0 \quad n \ge 2 \quad a_1(r) = 0.$$

for r = -p, we obtain

$$y_2(x) = x^{-p} \sum_{n \ge 0} \frac{(-1)^n}{2^{2n} n! (1-p) \dots (n-p)} x^{2n}$$

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Multiplying by $\frac{1}{2^{-p}\Gamma(1-p)}$ $J_{-p}(x) := \left(\frac{x}{2}\right)^{-p} \sum_{n \ge 0} \frac{(-1)^n}{n! \, \Gamma(n-p+1)} \, \left(\frac{x}{2}\right)^{2n} \quad x > 0.$

This is a second solution of the Bessel equation linearly independent of $J_p(x)$.

Case 1: 2p is not an integer.

Solving the recursion

$$[(r+n)^2 - p^2]a_n(r) + a_{n-2}(r) = 0 \quad n \ge 2 \quad a_1(r) = 0.$$

for r = -p, we obtain

$$y_2(x) = x^{-p} \sum_{n \ge 0} \frac{(-1)^n}{2^{2n} n! (1-p) \dots (n-p)} x^{2n}$$

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Case 2: 2p is a positive integer.

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Recall that the second solution is given by

$$y_2(x) = \sum_{n \ge 0} A'_n(-p)x^{n-p} + \sum_{n \ge 0} A_n(-p)x^{n-p}\log x$$

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$$a_{2n}(r) = \frac{(-1)^n}{\prod_{i=1}^n ((r+2i)^2 - p^2)}$$

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Thus, $A_{2n}(-p) = 0$ and $A'_{2n}(-p) = a_{2n}(-p)$. Thus, in this case we obtain that the second solution is

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$$J_{-p}(x) := \left(\frac{x}{2}\right)^{-p} \sum_{n \ge 0} \frac{(-1)^n}{n! \,\Gamma(n-p+1)} \,\left(\frac{x}{2}\right)^{2n} \quad x > 0.$$

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As before, $A_{2n+1}(r) = 0$. The polynomial $\prod_{i=1}^{n} ((r+2i)^2 - p^2)$ evaluated at r = -p, is $\prod_{i=1}^{n} 4i(i-p)$, Thus, if n < p, then $a_{2n}(r)$ is analytic in a neighborhood of -p. Thus, if n < p, then $A_{2n}(-p) = 0$ and

$$A'_{2n}(-p) = a_{2n}(-p) = \frac{(-1)^n}{2^{2n}n!(1-p)\dots(n-p)} = \frac{1}{2^{2n}n!(p-n)!}$$

If $n\geq p$, then

$$A_{2n}(-p) = \frac{2(-1)^n}{2^{2n}n!(1-p)\dots(-1)\cdot 1\cdot 2\cdots(n-p)}$$
$$= \frac{-2(-1)^{n-p}}{2^{2n}n!(p-1)!(n-p)!}$$

Define

$$f(r) := \Big(\prod_{i=1}^{p-1} ((r+2i)^2 - p^2)\Big)(r+3p)\Big(\prod_{i=p+1}^n ((r+2i)^2 - p^2)\Big) \quad (*)$$

Then

$$A_{2n}(r)f(r) = (-1)^n$$

Differentiating the above and setting r = -p we get

$$A'_{2n}(-p)f(-p) + A_{2n}(-p)f'(-p) = 0$$

Taking log and differentiating (\ast) we get

$$f'(-p) = f(-p) \left(\frac{1}{2p} + \sum_{i \in \{1, 2, \dots, n\} \setminus p} \frac{1}{2i} + \frac{1}{2(i-p)} \right)$$
$$= f(-p) \left(\frac{H_n}{2} - \frac{H_{p-1}}{2} + \frac{H_{n-p}}{2} \right),$$

where

$$H_0 = 0, \qquad H_n = 1 + \frac{1}{2} + \dots + \frac{1}{n}$$

Thus,

$$A'_{2n}(-p) = -A_{2n}(-p)\left(\frac{H_n}{2} - \frac{H_{p-1}}{2} + \frac{H_{n-p}}{2}\right)$$
$$= \frac{2(-1)^{n-p}}{2^{2n}n!(p-1)!(n-p)!}\left(\frac{H_n}{2} - \frac{H_{p-1}}{2} + \frac{H_{n-p}}{2}\right)$$

Thus, we get

$$y_{2}(x) = \sum_{n=0}^{p-1} \frac{1}{2^{2n} n! (p-n)!} x^{2n-p} + \sum_{n \ge p} \frac{(-1)^{n-p}}{2^{2n} n! (p-1)! (n-p)!} \Big(H_{n} - H_{p-1} + H_{n-p} \Big) x^{2n-p} + -\sum_{n \ge p} \frac{2(-1)^{n-p}}{2^{2n} n! (p-1)! (n-p)!} x^{2n-p} \log x$$

is a second solution.

Case 3: p = 0 (Repeated root case)

$$a_{2n}(r) = \frac{(-1)^n}{(r+2)^2(r+4)^2\dots(r+2n)^2} \qquad a_{2n+1}(r) = 0$$

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Differentiating $a_{2n}(r)$ with respect to r, we get

$$a_{2n}(r)' = -2a_{2n}(r)\left(\frac{1}{r+2} + \frac{1}{r+4} + \dots + \frac{1}{r+2n}\right)$$

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By theorem stated earlier, the second solution is

$$y_2(x) = J_0(x) \ln x - \sum_{n \ge 1} \frac{(-1)^n H_n}{2^{2n} (n!)^2} x^{2n} \quad x > 0$$

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$$y_2(x) = J_0(x) \ln x - \sum_{n \ge 1} \frac{(-1)^n H_n}{2^{2n} (n!)^2} x^{2n} \quad x > 0$$

where $y_1(x) = J_0(x) = \sum_{n \ge 0} \frac{(-1)^n}{2^{2n} (n!)^2} x^{2n}$ is Frobenius solution.

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Summary of p = 0 and p = 1/2

For p = 0, two independent solutions are $J_0(x)$, which is a real analytic function for all \mathbb{R} , and

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$$J_{\frac{1}{2}}(x) = \sqrt{\frac{2}{\pi x}} \sin x \text{ and } J_{-\frac{1}{2}}(x) = \sqrt{\frac{2}{\pi x}} \cos x$$

Both exhibit singular behavior at 0. Near 0, $J_{1/2}(x)$ is bounded but does not have a bounded derivative, while $J_{-1/2}(x)$ is unbounded near 0.

For real p, define

$$J_p(x) := \sum_{n=0}^{\infty} \frac{(-1)^n}{n! \, \Gamma(p+n+1)} \left(\frac{x}{2}\right)^{2n+p}$$

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- The above is a well defined power series once we know that the Gamma function never vanishes.
- If p ∉ {0, 1, 2, ...} J_p(x) and J_{-p}(x) are the two independent solutions of the Bessel equation.
- If $p \in \{0, 1, 2, ...\}$ then $J_{-p}(x) = (-1)^p J_p(x)$. Thus, in this case the second solution is not $J_{-p}(x)$.

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•
$$\frac{d}{dx} [x^p J_p(x)] = x^p J_{p-1}(x)$$

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These follow from (1) and (2). Expand LHS and divide by
$$x^{\pm p}$$
;

7

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$$J_{p-1}(x) - J_{p+1}(x) = 2J_p(x)$$

• $J_{p-1}(x) + J_{p+1}(x) = \frac{2p}{x}J_p(x)$

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The above two can be obtained by formally differentiating the power series.

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Problem: Let p > 0. Show that between any two <u>consecutive</u> zeros of $J_p(x)$, there exists <u>precisely one</u> zero of $J_{p-1}(x)$ and precisely one zero of $J_{p+1}(x)$

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Solution: Let 0 < c < d be two consecutive zeros of $J_p(x)$.

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 $[x^p J_p(x)]'(b) = 0 \quad \text{ for some } b \in (c, d)$

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$$[x^{p}J_{p}(x)]' = x^{p}J_{p-1}(x)$$

we get $J_{p-1}(b) = 0$. Repeating the above argument with the identity $[x^{-p}J_p(x)]' = -x^{-p}J_{p+1}(x)$, we get that $J_{p+1}(x)$ has a root in (c, d).

If $J_{p-1}(x)$ had two roots in (c, d), then from above, we conclude that $J_p(x)$ would have a root in (c, d). However, this contradicts the assumption that c and d are consecutive roots. Thus, J_{p-1} has exactly one root in (c, d).

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Solution: Using $J_{p-1}(x) - J_{p+1}(x) = 2J'_p(x)$ for p = 1, we get

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Now using $[x^{-p}J_p(x)]' = -x^{-p}J_{p+1}$ for p = 0, we get

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Therefore, $J_2(x) - J_0(x) = -2J'_1(x) = 2J''_0(x).$

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Therefore, $J_2(x) - J_0(x) = -2J'_1(x) = 2J''_0(x).$

Hence a = 2 and c = 0.

$$J_{p-1}(x) + J_{p+1}(x) = \frac{2p}{x} J_p(x)$$

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to see that $J_p(x)$ are elementary functions for $p\in \mathbb{Z}+\frac{1}{2}.$

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For example,

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$$J_{3/2}(x) = \frac{1}{x} J_{1/2}(x) - J_{-1/2}(x)$$

= $\sqrt{\frac{2}{\pi x}} \left(\frac{\sin x}{x} - \cos x \right)$

$$J_{p-1}(x) + J_{p+1}(x) = \frac{2p}{x} J_p(x)$$
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= $\sqrt{\frac{2}{\pi x}} \left(\frac{\sin x}{x} - \cos x\right)$

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$$J_{\frac{5}{2}}(x) = \frac{3}{x} J_{\frac{3}{2}}(x) - J_{\frac{1}{2}}(x)$$

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These functions are called spherical Bessel functions as they arise in solving wave equations in spherical coordinates.

Theorem (Liouville)

 $J_{m+\frac{1}{2}}(x)\, {\rm 's}$ are the only elementary Bessel functions.

Remark. Integrating some of the Bessel identities we get

$$\frac{d}{dx} [x^p J_p(x)] = x^p J_{p-1}(x)$$
$$\implies \int_0^x t^p J_{p-1}(t) dt = x^p J_p(x) + c$$

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For example,

$$\int_{0}^{x} t J_{0}(t) \, dt = x J_{1}(x) + c$$

(ロ)、(部)、(言)、(言)、(言)、(の)、(の) 27/56 It is rarely possible to solve 2nd order linear ODE

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Then how do we understand the nature and properties of solutions.

It is surprising that we can obtain quite a bit of information about the solution from the ODE itself.

If $y_1(x)$ and $y_2(x)$ are linearly independent solns of

$$y'' + P(x)y' + Q(x)y = 0$$

P, Q continuous on (a, b). Then (1) $y_1(x)$ and $y_2(x)$ have no common zero in (a, b). (2) Between any two successive zeros of $y_1(x)$, there is exactly one zero of $y_2(x)$ and vice versa.

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Proof of (1). Consider the Wronskian

$$W(x) := W(y_1, y_2) = y_1(x)y_2'(x) - y_1'(x)y_2(x)$$

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$$W(x) = C \exp\left(\int_{a_0}^x -P(t)dt\right) \qquad a_0 \in (a,b)$$

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In particular, since y_1 and y_2 are linearly independent, the Wronskian is nonzero and so it never vanishes. This proves (1).

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 $0 \neq W(x_1) = -y'_1(x_1)y_2(x_1) \qquad \qquad 0 \neq W(x_2) = -y'_1(x_2)y_2(x_2)$

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If $y_2(x)$ had two zeros in the interval $x_1 < \alpha < \beta < x_2$, then by the same reasoning, y_1 will have a zero in (α, β) , which contradicts the assumption that x_1 and x_2 are successive zeros of y_1 .

In particular, either both have finite number of zeros or both have infinite number of zeros in (a, b).

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One easily checks that u(x) satisfies the differential equation

$$u'' + q(x)u = 0 \qquad q(x) := Q(x) - \frac{1}{4}P(x)^2 - \frac{1}{2}P'(x)$$

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It is clear that the zeros of u are the same as those of y.

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It is clear that the zeros of u are the same as those of y. Also note that we need P(x) to be once differentiable.

Let u(x) be a non-trivial solution of u'' + q(x)u = 0 on the interval (α, β) , with q(x) continuous. Let $[a, b] \subset (\alpha, \beta)$ be a finite interval. Then u(x) has at most finite number of zeros in [a, b]. Hence if u(x) has infinitely many zeros on $(0, \infty)$, then the set of zeros of u(x) are not bounded.

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Proof. Assume u(x) has infinitely many zeros in [a, b]. Then $\exists x_0 \in [a, b]$ and a sequence of zeros $x_n \neq x_0$ such that $x_n \to x_0$ as $n \to \infty$.

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$$u'(x_0) = \lim_{x_n \to x_0} \frac{u(x_n) - u(x_0)}{x_n - x_0} = 0$$

This contradicts the fact that the Wronskian at x_0 is nonzero.

Let u(x) be a non-trivial solution of u'' + q(x)u = 0. If q(x) < 0 in (a,b) and continuous then u(x) has <u>atmost one zero</u> in (a,b).

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Then u(x) > 0 on (x_0, x_1) .

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Since u''(x) = -q(x)u(x) > 0 on (x_0, x_1) , u'(x) is an increasing function on (x_0, x_1) .

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But this is not possible as u' is increasing on (x_0, x_1) .

Let u(x) be a non-trivial solution of u'' + q(x)u = 0 Let q(x) be continuous and q(x) > 0 for all $x > x_0 > 0$. If $\int_{x_0}^{\infty} q(x) dx = \infty$, then u(x) has infinitely many zeros on $(0, \infty)$.

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$$u(x) = \int_{x_2}^x u'(t)dt + u(x_2) \le \int_{x_2}^x u'(x_2)dt + u(x_2)$$
$$\le u'(x_2)(x - x_2) + u(x_2)$$

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$$\begin{split} v(x) &= -\frac{u'(x)}{u(x)}, \quad \text{for} \ \ x \geq x_1 \\ v' &= \frac{-u''u + u'^2}{u^2} = \frac{q(x)u^2 + u'^2}{u^2} = q(x) + v(x)^2 \end{split}$$

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Integrating we get

$$v(x) - v(x_1) = \int_{x_1}^x q(x) \, dx + \int_{x_1}^x v(x)^2 \, dx$$

 $\int_{x_0}^{\infty} q(x) \, dx = \infty \implies v(x) > 0 \text{ for large } x.$

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 $\int_{x_0}^{\infty} q(x) \, dx = \infty \implies v(x) > 0 \text{ for large } x.$ Thus, u'(x) = -u(x)v(x) and this shows that u'(x) < 0 for x large.

In Bessel equation $x^2y'' + xy' + (x^2 - p^2)y = 0$ Substituting $u(x) = \sqrt{x}y(x)$, we get

$$u'' + \left[1 + \frac{1 - 4p^2}{4x^2}\right]u = 0$$

 $q(x) = 1 + \frac{1 - 4p^2}{4x^2}$ is continuous and q(x) > 0 for $x > x_0 > 0$. Further, $\int_{-\infty}^{\infty} \left(1 + \frac{1 - 4p^2}{4x^2}\right) dx = \infty$

$$\int_{x_0} \left(1 + \frac{1 - 4p}{4x^2} \right) \, dx = \infty$$

By previous theorem, u(x), hence any Bessel function has infinitely many zeros on $(0, \infty)$.

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Write $Z^{(p)} = \{x_1, x_2, ...\}$ as increasing sequence $x_n < x_{n+1}$. Question. What is the limit of $x_{n+1} - x_n$ as $n \to \infty$?

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Theorem (Sturm Comparison theorem)

Let y(x) be a non-trivial solutions of

y'' + q(x)y = 0

and z(x) be a non-trivial solutions of

z'' + r(x)z = 0

where q(x) > r(x) > 0 are continuous. Then y(x) vanishes at least once between any two consecutive zeros of z(x). Theorem (Sturm Comparison theorem)

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Compare y'' + 4y = 0 and z'' + z = 0. Here (q(x) =) 4 > (r(x) =) 1 > 0 Theorem (Sturm Comparison theorem)

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Here $(q(x) =) 4 > (r(x) =) 1 > 0$

Zeros of y(x) are $\pi/2$ apart and that of z(x) are π apart.

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Integrating from x_1 to x_2 , we get

$$W(x_2) - W(x_1) > 0 \implies W(x_2) > W(x_1)$$

But $W(x_1) = y(x_1)z'(x_1) > 0$ and $W(x_2) = y(x_2)z'(x_2) < 0$, a contradiction.

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• p < 1/2 \implies Between any two roots of $\alpha \cos x + \beta \sin x$ there is a root of $y_p(x)$.

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To see this, consider the function f := u(x - b + a) defined on the interval (b, ∞) .

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To see this, consider the function f := u(x - b + a) defined on the interval (b, ∞) .

It is a trivial check that f satisfies the differential equation

$$f'' + r(x)f = 0$$
 $r(x) := q(x - b + a)$

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Next we claim that the difference between any two successive roots of u is strictly less than π .

If not, then let a < b be successive roots such that $b-a \geq \pi$

If not, then let a < b be successive roots such that $b - a \ge \pi$ Since u has infinitely many roots, and their difference is strictly increasing, we may assume that $b - a > \pi$. If not, then let a < b be successive roots such that $b - a \ge \pi$ Since u has infinitely many roots, and their difference is strictly increasing, we may assume that $b - a > \pi$.

But now we can choose $\alpha, \beta \in \mathbb{R}$ such that $\alpha \cos x + \beta \sin x$ has two roots in (a, b), which contradicts Sturm's comparison theorem.

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Thus, we have proved that if $\{x_n\}$ are the roots of u in increasing order, then the difference $x_{n+1} - x_n$ is strictly increasing and bounded above by π .
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Thus, we have proved that if $\{x_n\}$ are the roots of u in increasing order, then the difference $x_{n+1} - x_n$ is strictly increasing and bounded above by π .

Next let us show that these differences converge to π . If not, then $(x_{n+1} - x_n) \rightarrow \gamma < \pi$. Choose $1 < \delta$, sufficiently close to 1 such that $\gamma < \frac{\pi}{\delta} < \pi$.

The function q(x) is decreasing to 1. Therefore, there is a $x_0 \in \mathbb{R}$, sufficiently large, such that $q(x_0) < \delta^2$. Apply Sturm's comparison on the interval (x_0, ∞) to the differential equations u'' + q(x)u = 0 and $z'' + \delta^2 z = 0$.

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Thus, between any two roots of u there is a root of z. Let a and b be two consecutive roots of u such that $x_0 < a < b$. Since $b - a < \gamma < \frac{\pi}{\delta}$, find a' and b' such that $x_0 < a' < a < b < b'$ and $b' - a' = \frac{\pi}{\delta}$.

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Find α and β such that the function $\alpha \cos \delta x + \beta \sin \delta x$ vanishes at a'. This function is a solution to the ODE $z'' + \delta^2 z = 0$. The next root of this function is at $a' + \frac{\pi}{\delta} = b'$. Thus, we get a contradiction to Sturm's theorem which says that there is a root of this function in the interval (a, b). Thus, we have proved

Theorem

If p < 1/2 then the sequence of differences of roots of u, $x_{n+1} - x_n$ is increasing and tends to π .

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If p < 1/2 then the sequence of differences of roots of u, $x_{n+1} - x_n$ is increasing and tends to π .

Similarly, we can prove that if p > 1/2 then the sequence of difference of roots of u is decreasing and tends to π .



The first few zeroes of Bessel functions are tabulated below.

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	$J_0(x)$	$J_1(x)$	$J_2(x)$	$J_3(x)$	$J_4(x)$	$J_5(x)$
1	2.4048	3.8317	5.1356	6.3802	7.5883	8.7715
2	5.5201	7.0156	8.4172	9.7610	11.0647	12.3386
3	8.6537	10.1735	11.6198	13.0152	14.3725	15.7002
4	11.7915	13.3237	14.7960	16.2235	17.6160	18.9801
5	14.9309	16.4706	17.9598	19.4094	20.8269	22.2178

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Question. Why are we concerned with zeros of Bessel function $J_p(x)$?

It is often required in mathematical physics to expand a given function in terms of Bessel functions.

Simplest and most useful expansions are of the form

$$f(x) = \sum_{n=1}^{\infty} a_n J_p(\lambda_{p,n} x) = a_1 J_p(\lambda_{p,1} x) + a_2 J_p(\lambda_{p,2} x) + \dots$$

where f(x) is defined on, (say) [0,1], and $\lambda_{p,n}$'s are zeros of Bessel function $J_p(x)$, $p \ge 0$.

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Remark: For a scalar a, the scaled Bessel functions $J_p(ax)$ are solutions of

$$x^2y'' + xy' + (a^2x^2 - p^2)y = 0$$

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known as scaled Bessel equation.

Orthogonality

 Define an inner product on functions on $\left[0,1\right]$ by

$$\langle f,g \rangle := \int_0^1 x f(x) g(x) \, dx$$

Define an inner product on functions on [0,1] by

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This is similar to the previous inner product except that f(x)g(x) is now multiplied by x and the interval of integration is from 0 to 1.

We call a function on $\left[0,1\right]$ square integrable with respect to this inner product if

$$\int_0^1 x f(x)^2 dx < \infty$$

The multiplying factor x is called a weight function.

Fix $p \ge 0$. Let $Z^{(p)} = \{\lambda_{p,1}, \lambda_{p,2}, \ldots\}$ denote the set of zeros of $J_p(x)$ on $(0, \infty)$.

Theorem

The set of scaled Bessel functions

 $\{J_p(\lambda_{p,1}x), J_p(\lambda_{p,2}x), \ldots\}$

form an orthogonal family w.r.t. above inner product, i.e. $\langle J_p(\lambda_{p,k}x), J_p(\lambda_{p,l}x) \rangle :=$

$$\int_0^1 x J_p(\lambda_{p,k} x) J_p(\lambda_{p,l} x) \, dx = \begin{cases} \frac{1}{2} [J_{p+1}(\lambda_{p,k})]^2 & \text{if } k = l \\ 0 & \text{if } k \neq l \end{cases}$$

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Theorem

Fix $p \ge 0$ and $Z^{(p)} = \{\lambda_{p,1}, \lambda_{p,2}, \ldots\}$: zeros of $J_p(x)$ on $(0, \infty)$. Any square-integrable function f(x) on [0,1] can be expanded in a series of scaled Bessel functions $J_p(\lambda_{p,n}x)$ as

$$f(x) = \sum_{n \ge 1} c_n J_p(\lambda_{p,n} x)$$

where

$$c_n = \frac{2}{[J_{p+1}(\lambda_{p,n})]^2} \int_0^1 x f(x) J_p(\lambda_{p,n} x) dx$$

This is Fourier-Bessel series of f(x) for parameter p.

Example. Let us compute the Fourier-Bessel series (for p = 0) of f(x) = 1 in the interval $0 \le x \le 1$.

$$\int_0^1 x \, J_0(\lambda_{0,n} x) \, dx = \frac{1}{\lambda_{0,n}} x \, J_1(\lambda_{0,n} x) \Big|_0^1 = \frac{J_1(\lambda_{0,n})}{\lambda_{0,n}}$$

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$$c_n = \frac{2}{[J_1(\lambda_{0,n})]^2} \int_0^1 x f(x) J_0(\lambda_{0,n}x) dx = \frac{2}{\lambda_{0,n} J_1(\lambda_{0,n})}$$

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Thus, the Fourier-Bessel series of f(x) is

$$\sum_{n\geq 1} \frac{2}{\lambda_{0,n} J_1(\lambda_{0,n})} J_0(\lambda_{0,n} x)$$

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By next theorem, this converges to 1 for 0 < x < 1.

Fourier-Bessel series converges to f(x) in norm, i.e.

$$\|f(x)-\sum_{n=1}^m c_n J_p(\lambda_{p,n} x)\| \ \text{ converges to } \ 0 \ \text{ as } \ m \to \infty$$

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For pointwise convergence, we have

Bessel expansion theorem

Assume f and f' have at most a finite number of jump discontinuities in [0,1], then the Bessel series converges for 0 < x < 1 to

$$\frac{f(x_-) + f(x_+)}{2}$$

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$$\frac{f(x_{-}) + f(x_{+})}{2}$$

At x = 1, the series always converges to 0 for all f, at x = 0, if p = 0 then it converges to $f(0_+)$. at x = 0, if p > 0 then it converges to 0. Proof of orthogonality of scaled Bessel functions

Proof of orthogonality of scaled Bessel functions If a, b are positive scalars, then $u(x) = J_p(ax)$ and $v(x) = J_p(bx)$ satisfies

$$u'' + \frac{1}{x}u' + \left(a^2 - \frac{p^2}{x^2}\right)u = 0$$
$$v'' + \frac{1}{x}v' + \left(b^2 - \frac{p^2}{x^2}\right)v = 0$$

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Multiply by v and u resp. and subtract, we get

$$(vu'' - uv'') + \frac{1}{x}(vu' - uv') + (a^2 - b^2)uv = 0$$

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Multiply by v and u resp. and subtract, we get

$$(vu'' - uv'') + \frac{1}{x}(vu' - uv') + (a^2 - b^2)uv = 0$$

$$(u'v - v'u)' + \frac{1}{x}(u'v - v'u) = (b^2 - a^2)uv$$

$$(x(u'v - v'u))' = (b^2 - a^2)xuv$$

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$$(b^{2} - a^{2}) \int_{0}^{1} xuv \, dx = \left[x(u'v - v'u) \right]_{0}^{1} = (u'v - v'u)(1)$$

$$(b^{2} - a^{2}) \int_{0}^{1} xuv \, dx = \left[x(u'v - v'u) \right] \Big|_{0}^{1} = (u'v - v'u)(1)$$
$$(b^{2} - a^{2}) \int_{0}^{1} xJ_{p}(ax)J_{p}(bx) \, dx = J'_{p}(a)J_{p}(b) - J'_{p}(b)J_{p}(a)$$

$$(b^{2} - a^{2}) \int_{0}^{1} xuv \, dx = \left[x(u'v - v'u) \right]_{0}^{1} = (u'v - v'u)(1)$$

$$(b^{2} - a^{2}) \int_{0}^{1} x J_{p}(ax) J_{p}(bx) \, dx = J_{p}'(a) J_{p}(b) - J_{p}'(b) J_{p}(a)$$

So if $a = \lambda_{p,k}$ and $b = \lambda_{p,l}$ are **distinct**, then

$$\int_0^1 x J_p(\lambda_{p,k} x) J_p(\lambda_{p,l} x) \, dx = 0$$

$$(b^{2} - a^{2}) \int_{0}^{1} xuv \, dx = \left[x(u'v - v'u) \right]_{0}^{1} = (u'v - v'u)(1)$$

$$(b^{2} - a^{2}) \int_{0}^{1} x J_{p}(ax) J_{p}(bx) \, dx = J_{p}'(a) J_{p}(b) - J_{p}'(b) J_{p}(a)$$

So if $a = \lambda_{p,k}$ and $b = \lambda_{p,l}$ are **distinct**, then

$$\int_0^1 x J_p(\lambda_{p,k} x) J_p(\lambda_{p,l} x) \, dx = 0$$

To compute the norm of $J_p(\lambda_{p,k}x)$, consider

$$2x^{2}u'\left[u'' + \frac{1}{x}u' + (a^{2} - \frac{p^{2}}{x^{2}})u\right] = 0$$
$$= [x^{2}u'^{2} + (a^{2}x^{2} - p^{2})u^{2}]' - 2a^{2}xu^{2}$$

Integrate on
$$[0,1]$$
,
$$2a^2\int_0^1xu^2\,dx=[x^2u'^2+(a^2x^2-p^2)u^2]\Big|_0^1$$

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Integrate on
$$[0,1]$$
,
$$2a^2\int_0^1 xu^2\,dx = [x^2u'^2 + (a^2x^2 - p^2)u^2]\Big|_0^1$$
Since $p \ge 0$, $(pu(0))^2 = (pJ_p(0))^2 = 0$.
Thus, $\left(x^2u'^2 + (a^2x^2 - p^2)u^2\right)(0) = 0$.
$$\begin{split} \text{Integrate on } [0,1], & 2a^2 \int_0^1 x u^2 \, dx = \left[x^2 u'^2 + (a^2 x^2 - p^2) u^2 \right] \Big|_0^1 \\ \text{Since } p \geq 0, \ (pu(0))^2 = (p J_p(0))^2 = 0. \\ \text{Thus, } \left(x^2 u'^2 + (a^2 x^2 - p^2) u^2 \right) (0) = 0. \\ \text{Further, } u'(1) = a J'_p(a), \text{ so we get} \\ & \left(x^2 u'^2 + (a^2 x^2 - p^2) u^2 \right) (1) = a^2 J'_p(a)^2 + (a^2 - p^2) J_p(a)^2 \end{split}$$

$$\begin{split} \text{Integrate on } [0,1], \\ & 2a^2 \int_0^1 x u^2 \, dx = \left[x^2 u'^2 + (a^2 x^2 - p^2) u^2 \right] \Big|_0^1 \\ \text{Since } p \geq 0, \ (pu(0))^2 = (p J_p(0))^2 = 0. \\ \text{Thus, } \ \left(x^2 u'^2 + (a^2 x^2 - p^2) u^2 \right) (0) = 0. \\ \text{Further, } u'(1) = a J'_p(a), \text{ so we get} \\ & \left(x^2 u'^2 + (a^2 x^2 - p^2) u^2 \right) (1) = a^2 J'_p(a)^2 + (a^2 - p^2) J_p(a)^2 \end{split}$$

Put $a = \lambda_{p,k}$ to get

$$2\lambda_{p,k}^2 \int_0^1 x J_p(\lambda_{p,k}x)^2 \, dx = \lambda_{p,k}^2 J_p'(\lambda_{p,k})^2$$

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$$\begin{split} \text{Integrate on } [0,1], \\ & 2a^2 \int_0^1 x u^2 \, dx = \left[x^2 u'^2 + (a^2 x^2 - p^2) u^2 \right] \Big|_0^1 \\ \text{Since } p \geq 0, \ (pu(0))^2 = (p J_p(0))^2 = 0. \\ \text{Thus, } \ \left(x^2 u'^2 + (a^2 x^2 - p^2) u^2 \right) (0) = 0. \\ \text{Further, } u'(1) = a J'_p(a), \text{ so we get} \\ & \left(x^2 u'^2 + (a^2 x^2 - p^2) u^2 \right) (1) = a^2 J'_p(a)^2 + (a^2 - p^2) J_p(a)^2 \end{split}$$

Put
$$a=\lambda_{p,k}$$
 to get
$$2\lambda_{p,k}^2\int_0^1x J_p(\lambda_{p,k}x)^2\,dx=\lambda_{p,k}^2J_p'(\lambda_{p,k})^2$$

Thus,

$$\int_0^1 x J_p(\lambda_{p,k}x)^2 \, dx = \frac{1}{2} J_p'(\lambda_{p,k})^2 = \frac{1}{2} J_{p+1}(\lambda_{p,k})^2$$

for last equality, use $x = \lambda_{p,k}$ in $J'_p(x) - \frac{p}{x}J_p(x) = J_{p+1}(x)$