#### MA-207 Differential Equations II

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For example, we will work out the case where

$$u(x,0) = f(x)$$
  $u(x,b) = 0$   $0 \le x \le a$   
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Let us apply the method of separation of variables. Let u(x,y)=A(x)B(y). Then the differential equation becomes

$$A''(x)B(y) + A(x)B''(y) = 0$$



Thus, we have

$$\frac{-A''(x)}{A(x)} = \frac{B''(y)}{B(y)} = \text{constant}$$

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For every  $n \ge 1$ , let

$$\lambda_n = \frac{n\pi}{a}$$

For each  $n \ge 1$ , we have a solution to

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Since we do not want A(x) to be identically 0 and u(x,b)=A(x)B(b)=0, this forces that Y(b)=0. Let us also impose the condition that Y(0)=1.

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Next consider for each  $\lambda_n$  the problem

$$B''(y) - \lambda_n^2 B(y) = 0$$
$$B(0) = 1$$
$$B(b) = 0$$

The solutions to the above equation are given by

$$B_n(y) = \sinh\left(\frac{n\pi(b-y)}{a}\right) / \sinh\left(\frac{n\pi b}{a}\right).$$

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$$u(x,y) = \sum_{n>1} \alpha_n \sin\left(\frac{n\pi x}{a}\right) \sinh\left(\frac{n\pi (b-y)}{a}\right) / \sinh\left(\frac{n\pi b}{a}\right),$$

where  $\alpha_n$  are real numbers.

#### Dirichlet boundary conditions: Formal solutions

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Thus, if f(x) has the Fourier expansion

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then we will have solved our Laplace equation with the given boundary conditions.

### Dirichlet boundary conditions: Formal solutions

#### Definition

Consider the Laplace equation with the boundary conditions

$$u_{xx} + u_{yy} = 0$$
  $0 < x < a, 0 < y < b$   
 $u(0, y) = 0 = u(a, y) = 0$   $0 \le y \le b$   
 $u(x, 0) = f(x)$   $0 \le x \le a$   
 $u(x, b) = 0$ 

The formal solution of the above problem is

$$u(x,t) = \sum_{n>1} \alpha_n \sin\left(\frac{n\pi x}{a}\right) \sinh\left(\frac{n\pi (b-y)}{a}\right) / \sinh\left(\frac{n\pi b}{a}\right),$$

where

$$\alpha_n = \frac{2}{L} \int_0^L f(x) \sin \frac{n\pi x}{L} \, dx$$

### Dirichlet boundary conditions: Actual solution

#### Theorem

Let f be continuous and piecewise smooth on [0,a] such that f(0)=f(a)=0. Consider the Laplace equation with the boundary conditions

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# Dirichlet boundary condition: Example

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$$u_{xx} + u_{yy} = 0 \qquad 0 < x < a, \quad 0 < y < b$$

$$u(0, y) = 0 = u(a, y) = 0 \qquad 0 \le y \le b$$

$$u(x, 0) = \sin\left(\frac{5\pi x}{a}\right) - 3\sin\left(\frac{9\pi x}{a}\right) \qquad 0 \le x \le a$$

$$u(x, b) = 0$$

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Since f is given by its Fourier series in the above example, it is clear that

$$\alpha_5 = 1$$

$$\alpha_9 = -3$$

# Dirichlet boundary condition: Example

#### Example (continued)

Thus, the solution to the above problem is given by

$$u(x,t) = \sin\left(\frac{5\pi x}{a}\right) \sinh\left(\frac{5\pi(b-y)}{a}\right) / \sinh\left(\frac{5\pi b}{a}\right)$$
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#### Neumann boundary condition

Consider the following differential equation

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For n > 0

$$B_0(y) = \frac{-1}{h}y + 1$$

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The formal solution of the above problem is

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$$\alpha_0 = \frac{1}{L} \int_0^L f(x) dx$$
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#### $\mathsf{Theorem}$

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Consider the Dirichlet problem in a disc of radius r

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on the boundary of the disc, which is a circle of radius r. To solve this problem write the Laplace operator in polar coordinates.

$$\Delta = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} = \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2}$$

#### Example

Solve for harmonic function  $u(r,\theta)$  in unit disc i.e.

$$\Delta u = 0, \quad r < 1, \ \theta \in [0, 2\pi]$$

with boundary condition

$$u(1,\theta) = f(\theta) = \begin{cases} \sin \theta, & \theta \in [0,\pi] \\ 0, & \theta \in [\pi, 2\pi] \end{cases}$$

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Laplace equation in polar coordinates is

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Again we try to analyze using separation of variables. Assume  $u(r,\theta)=R(r)\Theta(\theta).$  Then  $\Delta u=0$  becomes

$$R''(r)\Theta(\theta) + \frac{1}{r}R'(r)\Theta(\theta) + \frac{1}{r^2}R(r)\Theta''(\theta) = 0$$

From this we get

$$\frac{R''(r) + \frac{1}{r}R'(r)}{\frac{1}{r^2}R(r)} = -\frac{\Theta''(\theta)}{\Theta(\theta)} = \lambda.$$

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Thus, we need to solve

$$\Theta''(\theta) + \lambda\Theta(\theta) = 0$$
,  $r^2R''(r) + rR'(r) - \lambda R(r) = 0$ 

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Since  $u(r,\theta+2\pi)=u(r,\theta)$ , the functions  $\Theta$  and  $\Theta'$  need to be  $2\pi$  periodic.

From this we get

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Since  $u(r, \theta + 2\pi) = u(r, \theta)$ , the functions  $\Theta$  and  $\Theta'$  need to be  $2\pi$  periodic. Thus for the ODE for  $\Theta$ , we need to solve

$$\Theta''(\theta) + \lambda \Theta(\theta) = 0, \quad \Theta(0) = \Theta(2\pi), \quad \Theta'(0) = \Theta'(2\pi)$$

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Recall that this is EVP 5. The eigenvalues and eigenfunctions for periodic eigenvalue problem in  $\Theta$  are

$$\lambda_0=0,\;\;\Theta_0=1$$

and for  $n \ge 1$ ,

$$\lambda_n = n^2$$
,  $\Theta_{n,1}(\theta) = \cos(n\theta)$ ,  $\Theta_{n,2}(\theta) = \sin(n\theta)$ 

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The problem for R-function, namely

$$r^2R''(r) + rR'(r) - \lambda R(r) = 0$$
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is the Cauchy-Euler equation with solution  $x^m$ , where

$$m(m-1) + m - \lambda = m^2 - \lambda = 0.$$

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$$R_{0,1}(r) = 1$$
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Thus, the series solution has the form

$$u(r,\theta) = A_0 + \sum_{n>1} \left( A_n r^n \cos(n\theta) + B_n r^n \sin(n\theta) \right)$$

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Check that the Fourier series of  $f(\theta)$  is

$$f(\theta) = \frac{1}{\pi} - \frac{2}{\pi} \sum_{n \ge 1} \frac{\cos(2n\theta)}{4n^2 - 1} + \frac{1}{2}\sin\theta$$

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Therefore, the solution is

$$u(r,\theta) = \frac{1}{\pi} - \frac{2}{\pi} \sum_{n>1} \frac{1}{4n^2 - 1} r^{2n} \cos(2n\theta) + \frac{1}{2} r \sin\theta$$



#### Example

Solve for harmonic function  $u(r,\theta)$  in an annulus

$$\Delta u(r,\theta) = 0, \quad 1 < r < 2, \ \theta \in [0, 2\pi]$$
$$u(1,\theta) = \cos \theta, \quad 0 \le \theta \le 2\pi$$
$$u_r(2,\theta) = \sin 2\theta, \quad 0 \le \theta \le 2\pi$$

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This BVP can be interpreted as that for the steady state temperature distribution in an annular region where on the outer boundary the heat flux is prescribed and on the inner boundary, the temperature is prescribed.

Recall that the Laplace equation in polar coordinates is

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Again we try to analyze using separation of variables. Assume  $u(r,\theta)=R(r)\Theta(\theta).$  Then

$$R''(r)\Theta(\theta) + \frac{1}{r}R'(r)\Theta(\theta) + \frac{1}{r^2}R(r)\Theta''(\theta) = 0$$

From this we get

$$\frac{R''(r) + \frac{1}{r}R'(r)}{\frac{1}{r^2}R(r)} = -\frac{\Theta''(\theta)}{\Theta(\theta)} = \lambda$$

Thus, we need to solve

$$\Theta''(\theta) + \lambda \Theta(\theta) = 0, \quad r^2 R''(r) + rR'(r) - \lambda R(r) = 0$$

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Recall that this is EVP 5. The eigenvalues and eigen functions for periodic eigenvalue problem in  $\Theta$  are

$$\lambda_0 = 0, \quad \Theta_0 = 1$$

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The problem for R-function, namely

$$r^2R''(r) + rR'(r) - \lambda R(r) = 0$$

is the Cauchy-Euler equation with solution  $x^m$ , where

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Hence the general solution is

$$u(r,\theta) = (A_0 + B_0 \ln r) + \sum_{n \ge 1} (A_n r^n \cos(n\theta) + B_n n r^{-n} \cos(n\theta))$$
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We are given

$$u(1,\theta)=\cos\theta, \quad u_r(2,\theta)=\sin2\theta$$

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$$u_r(r,\theta) = \frac{B_0}{r} + \sum_{n \ge 1} n(A_n r^{n-1} - B_n r^{-n-1}) \cos n\theta$$
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$$+ n(C_n r^{n-1} - D_n r^{-n-1}) \sin n\theta$$

Comparing with  $u_r(2,\theta) = \sin 2\theta$ , we get  $B_0 = 0$ ,  $2(2C_2 - 2^{-3}D_2) = 1$ 

$$A_n 2^{n-1} - B_n 2^{-n-1} = 0 \ (n \ge 1), \quad C_n 2^{n-1} - D_n 2^{-n-1} = 0 \ (n \ne 2)$$

Using these equations we can solve for the values of the constants.

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,  $A_1 - B_1 2^{-2} = 0 \implies A_1 = \frac{1}{5}$ ,  $B_1 = \frac{4}{5}$ 

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For n=2,

$$A_2 + B_2 = 0$$
,  $A_2 - B_2 - B_2 = 0 \implies A_2 = 0 = B_2$ 

$$C_2 + D_2 = 0$$
,  $2C_2 - \frac{1}{2^3}D_2 = \frac{1}{2} \implies C_2 = \frac{4}{17}$ ,  $D_2 = \frac{-4}{17}$ 

For n > 2,

$$A_n + B_n = 0$$
,  $A_n 2^{n-1} - B_n 2^{-n-1} = 0 \implies A_n^1 = 0 = B_n^1$   
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Thus the solution is

$$u(r,\theta) = (\frac{1}{5}r + \frac{4}{5}r^{-1})\cos\theta + (\frac{4}{17}r^2 + \frac{-4}{17}r^{-2})\sin 2\theta$$

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Let use the method of separation of variables to get some solutions.

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Let  $u(r, \theta, t) = X(r)Y(\theta)T(t)$ .

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$$\frac{T''(t)}{k^2T(t)} = \frac{r^2X''(r) + rX'(r)}{r^2X(r)} - \frac{n^2}{r^2}$$

$$Y''(\theta) + n^2 Y(\theta) = 0$$

If n = 0, then the solution is  $Y(\theta) = \text{constant}$ .

If  $n \ge 1$ , then the solution is of the type

$$Y(\theta) = A\cos(n\theta) + B\sin(n\theta).$$

The second equation becomes

$$\frac{r^2T''(t)}{k^2T(t)} - \frac{r^2X''(r) + rX'(r)}{X(r)} = -n^2$$

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Again, this forces both sides to be a constant c.



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$$J_n(ar) = \left(\frac{ar}{2}\right)^n \sum_{m \ge 0} \frac{(-1)^m}{m! \Gamma(m+n+1)} \left(\frac{r}{2}\right)^{2m} \quad r > 0$$

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Clearly, this solution is unbounded as  $r \to 0$ .

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$$(A\cos(n\theta) + B\sin(n\theta))(C\cos(\mu_{n,i}kt) + D\sin(\mu_{n,i}kt))$$

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So 
$$u_{0,i}(r,\theta,t) = J_0(\mu_{0,i}r) \Big( C\cos(\mu_{n,i}kt) + D\sin(\mu_{n,i}kt) \Big)$$

$$(A\cos(n\theta) + B\sin(n\theta))(C\cos(\mu_{n,i}kt) + D\sin(\mu_{n,i}kt))$$

For n=0, we only allow constant solutions for Y (since we want the solutions to be periodic),

So 
$$u_{0,i}(r,\theta,t) = J_0(\mu_{0,i}r) \Big( C\cos(\mu_{n,i}kt) + D\sin(\mu_{n,i}kt) \Big)$$

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Thus, consider  $u(r,\theta,t) = \sum_{n \geq 0} \sum_{i \geq 1} \Big( A_{n,i}\cos(n\theta)\cos(\mu_{n,i}t) + A_{n,i}\cos(n\theta)\cos(\mu_{n,i}t) \Big)$ 

$$(A\cos(n\theta) + B\sin(n\theta))(C\cos(\mu_{n,i}kt) + D\sin(\mu_{n,i}kt))$$

So 
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Thus, consider 
$$u(r,\theta,t) = \sum_{n\geq 0,i\geq 1} \Big( A_{n,i}\cos(n\theta)\cos(\mu_{n,i}t) + \\ + B_{n,i}\sin(n\theta)\cos(\mu_{n,i}t) + C_{n,i}\cos(n\theta)\sin(\mu_{n,i}t) \\ + D_{n,i}\sin(n\theta)\sin(\mu_{n,i}t) \Big) J_n(\mu_{n,i}r)$$

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$$\sum_{n>0,i>1} (A_{n,i}\cos(n\theta) + B_{n,i}\sin(n\theta))J_n(\mu_{n,i}r)$$

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The initial condition that we have is  $u(r, \theta, 0) = f(r, \theta)$ Question: Can we write f as an expansion as above?

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$$S := \{J_n(\mu_{n,i}r)\cos(n\theta), J_m(\mu_{m,j}r)\sin(m\theta)\}_{n \ge 0, i \ge 1}$$

is an orthogonal set of functions, under the inner product

$$\langle f, g \rangle := \int_0^R \int_0^{2\pi} f(r, \theta) g(r, \theta) r d\theta dr$$

Next let us observe that for any function f(r) and  $n \ge 0$  fixed,

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This is because  $\{J_n(\mu_{n,i}r)\}_{i\geq 1}$  is a maximal orthogonal set for functions f on [0,R] such that  $\langle f,f\rangle<\infty$ .

Recall that for f as above,

where

$$c_i = \frac{1}{\langle J_n(\mu_{n,i}r), J_n(\mu_{n,i}r) \rangle}$$

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Similarly, the function  $f(r)\sin(n\theta)$  has a Fourier-Bessel series in terms of  $J_n(\mu_{n,i}r)\sin(n\theta)$ 

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If  $g(\theta)$  is any function in  $L^2([0,\theta])$ , then  $g(\theta)$  has a Fourier expansion in terms of  $\{\cos(n\theta),\sin(n\theta)\}_{n\in\mathbb{Z}}$ , we get that any function of the type  $f(r)g(\theta)$  can be approximated by the functions in S.

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$$A_{n,i} = \frac{\langle h, J_n(\mu_{n,i}r)\cos(n\theta)\rangle}{\langle J_n(\mu_{n,i}r)\cos(n\theta), J_n(\mu_{n,i}r)\cos(n\theta)\rangle}$$

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Similarly, we get  $B_{n,i}$  by replacing  $\cos$  by  $\sin$  in the above.

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#### $\mathsf{Theorem}$

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### Theorem (continued ..)

The above differential equation with initial and boundary conditions has a solution given by

$$u(r, \theta, t) = \sum_{n \ge 0, i \ge 1} \left( A_{n,i} \cos(n\theta) \cos(\mu_{n,i} t) + B_{n,i} \sin(n\theta) \cos(\mu_{n,i} t) + C_{n,i} \cos(n\theta) \sin(\mu_{n,i} t) + C_{n,i} \sin(n\theta) \sin(\mu_{n,i} t) \right)$$

### Theorem (continued ..)

The above differential equation with initial and boundary conditions has a solution given by

$$u(r, \theta, t) = \sum_{n \ge 0, i \ge 1} \left( A_{n,i} \cos(n\theta) \cos(\mu_{n,i} t) + B_{n,i} \sin(n\theta) \cos(\mu_{n,i} t) + C_{n,i} \cos(n\theta) \sin(\mu_{n,i} t) + C_{n,i} \sin(n\theta) \sin(\mu_{n,i} t) \right)$$

where

$$A_{n,i} = \frac{\langle f, J_n(\mu_{n,i}r)\cos(n\theta)\rangle}{\langle J_n(\mu_{n,i}r)\cos(n\theta), J_n(\mu_{n,i}r)\cos(n\theta)\rangle}$$

### Theorem (continued ..)

The above differential equation with initial and boundary conditions has a solution given by

$$u(r, \theta, t) = \sum_{n \ge 0, i \ge 1} \left( A_{n,i} \cos(n\theta) \cos(\mu_{n,i} t) + B_{n,i} \sin(n\theta) \cos(\mu_{n,i} t) + C_{n,i} \cos(n\theta) \sin(\mu_{n,i} t) + C_{n,i} \sin(n\theta) \sin(\mu_{n,i} t) \right)$$

where

$$A_{n,i} = \frac{\langle f, J_n(\mu_{n,i}r)\cos(n\theta)\rangle}{\langle J_n(\mu_{n,i}r)\cos(n\theta), J_n(\mu_{n,i}r)\cos(n\theta)\rangle}$$

$$B_{n,i} = \frac{\langle f, J_n(\mu_{n,i}r)\sin(n\theta)\rangle}{\langle J_n(\mu_{n,i}r)\sin(n\theta), J_n(\mu_{n,i}r)\sin(n\theta)\rangle}$$

## Theorem (continued ..)

$$C_{n,i} = \frac{1}{\mu_{n,i}} \frac{\langle g, J_n(\mu_{n,i}r) \cos(n\theta) \rangle}{\langle J_n(\mu_{n,i}r) \cos(n\theta), J_n(\mu_{n,i}r) \cos(n\theta) \rangle}$$

$$D_{n,i} = \frac{1}{\mu_{n,i}} \frac{\langle g, J_n(\mu_{n,i}r) \sin(n\theta) \rangle}{\langle J_n(\mu_{n,i}r) \sin(n\theta), J_n(\mu_{n,i}r) \sin(n\theta) \rangle}$$